# Exponential PDEs in high dimensions

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P. Esposito May 16, 2024 Analysis of Geometric Singularities

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 $(M^2, g)$  closed Riemannian 2-manifold,  $\Delta_g$  Laplace-Beltrami operator,  $0 = \lambda_0 < \lambda_1 \le \lambda_2 \le \ldots$  eigenvalues of  $-\Delta_g$ 

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$$\log \frac{\det(-\Delta_{\hat{g}})}{\det(-\Delta_g)} = -\frac{1}{12\pi} \int_M (|\nabla u|_g^2 + 2K_g u) \ dv_g, \quad \hat{g} = e^{2u}g$$

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Application: compactness of isospectral domains/surfaces:

- B. Osgood, R. Phillips, P. Sarnak, JFA '88
- B. Osgood, R. Phillips, P. Sarnak, Ann. Math. '89

 $A_g$  conformally covariant: if  $\hat{g} = e^{2u}g$  then  $A_{\hat{g}}\psi = e^{-bu}A_g(e^{au}\psi)$ 

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$$\begin{split} A_g \text{ conformally covariant: if } \hat{g} &= e^{2u}g \text{ then } A_{\hat{g}}\psi = e^{-bu}A_g(e^{au}\psi)\\ \underline{\text{Branson-Orsted formula: }} (M^4,g) \text{ closed Riemannian 4-manifold,}\\ &\ker A_g = \{0\}, \, \hat{g} = e^{2u}g \Rightarrow\\ F_{A_g}[u] &= \log \frac{\det A_{\hat{g}}}{\det A_{\sigma}} = \gamma_1 I[u] + \gamma_2 II[u] + \gamma_3 III[u] \quad (\gamma_i \in \mathbb{R}) \end{split}$$

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Examples:

- conformal Laplacian  $L_g = -\Delta_g + \frac{(n-2)}{4(n-1)}R_g$
- Paneitz operator  $P_g = \Delta_g^2 \operatorname{div}(\frac{2}{3}R_gg 2Ric_g) \circ \nabla$
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$$\Rightarrow \gamma_2(L_g) = 6\gamma_3(L_g), \quad \gamma_2(\not\!\!D_g^2) = \frac{132}{7}\gamma_3(\not\!\!D_g^2)$$

 $W_g$  Weyl tensor of g,  $\hat{g} = e^{2u}g$ 

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Each functional  $\leftrightarrow$  a natural curvature condition:

$$\begin{split} I'(u) &= 0 & \Leftrightarrow & |W_{\hat{g}}|^2 = const.\\ II'(u) &= 0 & \Leftrightarrow & Q_{\hat{g}} = const.\\ III'(u) &= 0 & \Leftrightarrow & \Delta_{\hat{g}}R_{\hat{g}} = 0 \end{split}$$

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$$P_{g}u + 2Q_{g} = 2Q_{\hat{g}}e^{4u}, \quad Q_{g} = \frac{1}{12}(-\Delta_{g}R_{g} + R_{g}^{2} - 3|Ric_{g}|_{g}^{2})$$

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Gauss-Bonnet formula:  $4\pi^{2}\chi(M) = \int_{M}(\frac{|W_{g}|_{g}^{2}}{8} + Q_{g}) dv_{g}$ , where  
(*M*) Euler characteristic of *M*
  
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Analysis of Geometric Singularities

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<u>E-L eqn</u>:  $\mathcal{N}_g(u) + U_g = -k_{A_g} \frac{e^{4u}}{\int_M e^{4u} dv_g}$  where

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$$\mathcal{N}_g(u) = \frac{\gamma_2}{2} P_g u + 6\gamma_3 \Delta_g (\Delta_g u + |\nabla u|_g^2) - 12\gamma_3 \text{div} [(\Delta_g u + |\nabla u|_g^2) \nabla u] + 2\gamma_3 \text{div} (R_g \nabla u)$$

<u>Difficulty</u>:  $\mathcal{N}_g(u) = (\frac{\gamma_2}{2} + 6\gamma_3)\Delta_g^2 u - 12\gamma_3\Delta_{4,g}u + \dots$  is a quasi-linear operator of mixed orders

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Existence of extremals: [A. Chang, P. Yang, Ann. Math. '95]

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$$\gamma_2, \gamma_3 < 0$$
 and  $k_{A_g} < 8\pi^2(-\gamma_2)$ 

•  $\gamma_1 = \gamma_3 = 0$ :  $P_g \ge 0$  with ker $P_g = \mathbb{R}$  and  $k_{P_g} < 8\pi^2$ 

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[M. Gursky, CMP '99]  $k_{L_g} < 32\pi^2$  if  $R_g \ge 0$  except on  $(\mathbb{S}^4, g_0)$ 

In general  $k_{A_g} \leq 8\pi^2(-\gamma_2)$  fails (products of negatively-curved surfaces, hyperbolic manifolds)

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Conformal metrics with Q = const. found as saddle points of II in

• Z. Djadli, A. Malchiodi, Ann. Math. '08

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#### Theorem 1 (P.E., A. Malchiodi, JDG '24)

 $\frac{\gamma_2}{\gamma_3} \ge 6$ ,  $w_n$  blow-up sequence of  $\mathcal{N}_g(w_n) + U_g = \mu_n e^{4w_n}$  in MThen  $\int_M w_n dv_g \to -\infty$  and  $\mu_n e^{4w_n} \rightharpoonup 8\pi^2 \gamma_2 \sum \delta_{p_i}$ 

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We find conformal metrics with U = const. as saddle points of  $F_{A_g}$ :

#### Theorem 2 (P.E., A. Malchiodi, JDG '24)

 $\frac{\gamma_2}{\gamma_3} \ge 6$ , *M* compact manifold s.t.  $k_{A_g} \notin 8\pi^2(-\gamma_2)\mathbb{N}$ Then  $\exists \ \tilde{g} \in [g]$  with  $U_{\tilde{g}} = const$ .

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- T. Branson, A. Chang, P. Yang, CMP '92 ["unique" maximizer for  $F_{L_{g_0}}$  and  $F_{D_{g_0}^2}$ ]
- M. Gursky, CMP '97 ["unique" c.p. for  $F_{L_{g_0}}$  and  $F_{\mathcal{D}_{g_0}^2}$ ]
- M. Gursky, A. Malchiodi, CMP '12 [non-uniqueness for  $\gamma_2 < 0 < \gamma_3$ ]

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Via stereographic projection  $F'_{L_{g_0}}(w) = 0$  on  $(\mathbb{S}^4, g_0)$  equivalent to

 $3\Delta^2 W + 2\Delta(|\nabla W|^2) - 4\text{div} \left[(\Delta W + |\nabla W|^2)\nabla W\right] = 16\pi^2 e^{4W}$ 

in  $\mathbb{R}^4$  with  $W \sim -2\log|x|$  at infinity.

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Via stereographic projection  $F'_{L_{g_0}}(w) = 0$  on  $(\mathbb{S}^4, g_0)$  equivalent to

 $3\Delta^2 W + 2\Delta(|\nabla W|^2) - 4\text{div} \left[(\Delta W + |\nabla W|^2)\nabla W\right] = 16\pi^2 e^{4W}$ 

in  $\mathbb{R}^4$  with  $W \sim -2 \log |x|$  at infinity. Solutions are classified as translations and dilations of

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in  $\mathbb{R}^4$  with  $W\sim -2\log|x|$  at infinity. Solutions are classified as translations and dilations of

$$W = \log(\frac{2}{1+|x|^2})$$
  
Question:  $\int_{\mathbb{R}^4} e^{4W} < +\infty$  does it implies  $W \sim -2\log|x|$  at  $\infty$ ?

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Simplified problem: retain  $\Delta_{4,g}$  in  $\mathcal{N}_g$  and consider it in general dimensions  $n\geq 2$ 

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Arises also in

 $-\Delta_{n,g}u + |\nabla u|_g^{n-2}\operatorname{Ric}_g(\nabla u, \nabla u) = |\nabla u|_{\hat{g}}^{n-2}\operatorname{Ric}_{\hat{g}}(\nabla u, \nabla u)e^{nu}$ 

see

• S. Ma, J. Qing, Calc. Var '21 & Adv. Math. '22

## The *n*-Liouville equation

Consider the euclidean quasilinear PDE

$$-\Delta_n u = V e^u \text{ in } \Omega, \quad u = 0 \text{ on } \partial \Omega \tag{P}$$

Analysis of Geometric Singularities

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<u>Planar case n = 2</u>: arises in conformal geometry, statistical and fluid mechanics, Chern-Simons theories;

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<u>Planar case n = 2</u>: arises in conformal geometry, statistical and fluid mechanics, Chern-Simons theories; well studied on Euclidean domains or on closed Riemannian surfaces

- H. Brézis, F. Merle, Comm. PDE '91
- Y.Y. Li, I. Shafrir, Indiana Univ. Math. J. '94
- Y.Y. Li, Comm. Math. Phys. '99
- C.C. Chen, C.S. Lin, Comm. Pure Appl. Math. '02 & '03

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Question: n > 2?

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Question: n > 2? A concentration-compactness principle:

- X. Ren, J. Wei, J. Differential Equations '95
- J.A. Aguilar, I. Peral, Nonlinear Anal. '97

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#### Theorem 3 (P.E., F. Morlando, JMPA '15)

Let  $u_k$  be solutions of (P) with  $V_k$  satisfying (V) and  $\sup_{k \in \mathbb{N}} \int_{\Omega} e^{u_k} < +\infty \quad \& \quad \sup_{k \in \mathbb{N}} osc_{\partial\Omega} u_k < +\infty$ 

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If  $osc_{\partial\Omega}u_k = 0$ , (i)-(ii) do hold in  $\overline{\Omega}$  with  $S \subset \Omega$  in case (ii).

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<u>Notation</u>:  $c_n = n(\frac{n^2}{n-1})^{n-1}$ 

P. Esposito May 16, 2024

Analysis of Geometric Singularities

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For  $\lambda \notin c_n \omega_n \mathbb{N}$  Theorem 3 gives compactness for solutions of

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- W. Ding, J. Jost, J. Li, G. Wang, AIHP '99
- Z. Djadli, A. Malchiodi, Ann. Math. '08
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- Z. Djadli, A. Malchiodi, Ann. Math. '08
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<u>Alternative approaches</u>: via degree (blow-up analysis misses); via perturbative methods (difficult due to nonlinearity of  $\Delta_n$ )

P. Esposito May 16, 2024

Analysis of Geometric Singularities

Aim: classify solutions of

$$-\Delta_n U = e^U$$
 in  $\mathbb{R}^n$ ,  $\int_{\mathbb{R}^n} e^U < \infty$   $(P)_{\infty}$ 

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Analysis of Geometric Singularities

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Scaling and translation invariance  $\Rightarrow$  explicit solutions  $U_{\lambda,p}$ :

$$U_{\lambda,p}(x) = \log rac{c_n \lambda^n}{(1+\lambda^{rac{n}{n-1}}|x-p|^{rac{n}{n-1}})^n} \qquad \lambda>0, \ p\in \mathbb{R}^n$$

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Quantization:  $\int_{\mathbb{R}^n} e^{U_{\lambda,p}} = c_n \omega_n$ 

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#### Theorem 5 (P.E., AIHP '18)

Any solution U of  $(P)_{\infty}$  has the form  $U_{\lambda,p}$ . In particular  $\int_{\mathbb{R}^n} e^U = c_n \omega_n$ 

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Analysis of Geometric Singularities

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- J. Liouville, J. de Math. 1853 [via complex analysis]
- W. Chen, C. Li, Duke Math. J. '91 [via moving planes]

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Liouville approach: integrability & the Liouville theorem in  $\mathbb C$ 

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<u>Chen-Li approach</u>: integral representation of U to deduce logarithmic behavior of U at  $\infty$  in terms of  $\int_{\mathbb{R}^2} e^U$ 

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<u>Chen-Li approach</u>: integral representation of U to deduce logarithmic behavior of U at  $\infty$  in terms of  $\int_{\mathbb{R}^2} e^U$ &  $\int_{\mathbb{R}^2} e^U \ge 8\pi$  via an isoperimetric argument  $\Rightarrow$  enough decay to carry out a simple MP approach

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• no integral representation for a solution U of  $(P)_{\infty}$ 

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An alternative approach: via Pohozev identity in

- P.-L. Lions, Appl. Anal. '81
- S. Kesavan, F. Pacella, Appl. Anal. '94
- S. Chanillo, M. Kiessling, Geom. Funct. Anal. '95

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If U is a solution of  $(P)_{\infty} \Rightarrow$  the Kelvin transform  $\hat{U}$  satisfies

$$-\Delta_n \hat{U} = \frac{e^{\hat{U}}}{|x|^{2n}} \text{ in } \mathbb{R}^n \setminus \{0\}, \ \int_{\mathbb{R}^n} \frac{e^{\hat{U}}}{|x|^{2n}} < +\infty$$

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Analysis of Geometric Singularities

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The description of singularities in

- J. Serrin, Acta Math. '64 and '65
- S. Kichenassamy, L. Veron, Math. Ann. '86

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fails in the limiting situation  $F \in L^1 \Rightarrow -\Delta \hat{U} = \frac{e^{\hat{U}}}{|x|^{2n}} - \left(\int_{\mathbb{R}^n} e^{U}\right) \delta_0$ &  $\hat{U}$  log. behavior at 0

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$$-\Delta_n U = e^U - \gamma \delta_0 ext{ in } \mathbb{R}^n, \quad \int_{\mathbb{R}^n} e^U < +\infty$$

• P. E., Calc. Var. PDE '21 [if  $n \ge 2$ ]

• J. Prajapat, G. Tarantello, Proc. Edinburgh '01 [if  $\underline{n} = 2$ ]
Dropping  $\sup_{k \in \mathbb{N}} \operatorname{osc}_{\partial \Omega} u_k < +\infty$ , in general concentration masses satisfy  $\alpha_p \ge n^n \omega_n$ 

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If  $0 \leq V_k \rightarrow V$  in  $C_{loc}(\Omega)$ , then  $\alpha_p \geq c_n \omega_n$  thanks to mass quantization for the limiting problem

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In the two-dimensional case  $\alpha_p \in 8\pi\mathbb{N}$  is shown in

• Y.Y. Li, I. Shafrir, Indiana Univ. Math. J. '94

based on a Harnack inequality of  $\sup + \inf type$ 

• I. Shafrir, C.R.A.S. '92

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The main point comes from the "linear theory", see also F. Robert, Proc. Edinburgh '07:

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$$u(x) - \inf_{\Omega} u \ge c_1 \int_0^{\delta} \left[ \int_{B_t(x)} f \right]^{\frac{1}{n-1}} \frac{dt}{t}$$

Analysis of Geometric Singularities

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If 
$$c_1 = (n\omega_n)^{-\frac{1}{n-1}}$$
, set  $\mu_k = e^{-\frac{u_k(x_k)}{n}}$  with  $u_k(x_k) = \max_K u_k$ :  
 $u_k(x_k) - \inf_{\Omega} u_k \ge \left[\frac{1}{n\omega_n}\int_{B_{R\mu_k}(x_k)} V_k e^{u_k}\right]^{\frac{1}{n-1}}\log\frac{\delta}{R\mu_k}$ 

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$$u(x) - \inf_{\Omega} u \ge c_1 \int_0^{\delta} \left[ \int_{B_t(x)} f \right]^{\frac{1}{n-1}} \frac{dt}{t}$$

If 
$$c_1 = (n\omega_n)^{-\frac{1}{n-1}}$$
, set  $\mu_k = e^{-\frac{u_k(x_k)}{n}}$  with  $u_k(x_k) = \max_K u_k$ :  
 $u_k(x_k) - \inf_{\Omega} u_k \ge \left[\frac{1}{n\omega_n}\int_{B_{R\mu_k}(x_k)} V_k e^{u_k}\right]^{\frac{1}{n-1}}\log\frac{\delta}{R\mu_k}$   
 $\Rightarrow u_k(x_k) - \inf_{\Omega} u_k \ge \left(\frac{n}{n-1} - \delta\right)u_k(x_k) + C$ 

for all  $\delta$  small in view of  $\int_{\mathcal{B}_{R\mu_k}(x_k)} V_k e^{u_k} \sim c_n \omega_n$ 

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#### Theorem 6 (P.E., M. Lucia, preprint)

Given  $K \subset \Omega$  compact and  $C_1 < \frac{1}{n-1}$ , there exists  $C_2 > 0$  so that  $C_1 \max_K u_k + \inf_\Omega u_k \le C_2$ 

The constant  $c_1$  is not explicit but  $0 < c_1 \le (n\omega_n)^{-\frac{1}{n-1}}$ , see

• T. Kilpeläinen, J. Malý, Ann. SNS Pisa '92

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If  $f \ge 0$  is radial in  $B_{\delta}(x)$ , then by comparison

$$u(x) - \inf_{\Omega} u \ge u(x) - \inf_{B_{\delta}(x)} u \ge \int_{0}^{\delta} \left(\frac{1}{n\omega_{n}} \int_{B_{t}(x)} f\right)^{\frac{1}{n-1}} \frac{dt}{t}$$

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<u>n=2</u>: by Green's representation formula for all  $y \in B_{\delta}(x)$ 

$$u(y) - \inf_{B_{\delta}(x)} u \geq \int_{B_{\delta}(x)} \Big[ -\frac{1}{2\pi} \log |z-y| + H(z,y) \Big] f(z) dz$$

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$$\Rightarrow \quad u(x) - \inf_{\Omega} u \ge -\frac{1}{2\pi} \int_{B_{\delta}(x)} \log \frac{|z-x|}{\delta} f(z) dz$$
$$= -\frac{1}{2\pi} \int_{0}^{2\pi} d\theta \int_{0}^{\delta} t \log \frac{t}{\delta} f(t\theta + x) dt$$
$$= \frac{1}{2\pi} \int_{0}^{2\pi} d\theta \int_{0}^{\delta} \frac{dt}{t} \int_{0}^{t} rf(r\theta + x) dr = \frac{1}{2\pi} \int_{0}^{\delta} [\int_{B_{\delta}(x)} f] \frac{dt}{t}$$

Since in general  $c_1 < (n\omega_n)^{-\frac{1}{n-1}}$ , we need to fill the gap via a blow-up approach:

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#### General case

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• linear theory still implies finite mass for the limiting profiles

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$$\begin{cases} -\Delta_n w_k = V_k e^{u_k} \chi_{B_{R\mu_k}(x_k)} & \text{in } B_{\delta}(x_k) \\ w_k = 0 & \text{on } \partial B_{\delta}(x_k) \end{cases}$$

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• since  $V_k e^{u_k} \sim V(p) e^{\bigcup_{\mu_k^{-1}, x_k}}$  in  $B_{R\mu_k}(x_k)$  with  $p = \lim_{k \to +\infty} x_k$ , further compare  $w_k$  from below with the radial case where  $c_1 = (n\omega_n)^{-\frac{1}{n-1}}$ 

#### Quantization for mass concentration

By sup + inf inequalities one gets decay estimates on  $V_k e^{u_k}$ :

$$V_k e^{u_k} \leq C rac{\mu_k^lpha}{|x-x_k|^{n+lpha}} \qquad ext{in } B_{rac{d_k}{2}}(x_k) \setminus B_{R\mu_k}(x_k)$$

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By  $\int_{B_{R\mu_k}(x_k)} V_k e^{u_k} \sim c_n \omega_n$  and decay estimates one gets that  $\int_{B_{\frac{d_k}{2}}(x_k)} V_k e^{u_k} \sim c_n \omega_n$ 

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Following clusters by clusters, rather standard to show that

## Theorem 7 (P.E., M. Lucia, preprint) $lpha_{p} \in c_{n}\omega_{n}\mathbb{N}$

extending the two-dimensional result in

• Y.Y. Li, I. Shafrir, Indiana Univ. Math. J. '94

### Optimal sup + inf inequality

The decay exponent  $\alpha$  is in general  $\alpha < \frac{n}{n-1}$ . When blow-up is simple, is it possible to reach  $\alpha = \frac{n}{n-1}$ ? Equivalent to

$$V_k e^{u_k} \le C rac{\mu_k^{rac{n}{n-1}}}{|x-x_k|^{rac{n^2}{n-1}}}$$

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The answer is related to the following fundamental expansion

$$u_k - U_{\mu_k^{-1}, \mathsf{x}_k} = O(1)$$
 in  $B_\delta(p)$ 

and optimal constant  $C_1 = \frac{1}{n-1}$  in the sup + inf inequality, see

- D. Bartolucci, C.C. Chen, C.S. Lin, G. Tarantello, Comm. PDE '04
- H. Brézis, Y.Y. Li, I. Shafrir, JFA '93
- Y.Y. Li, Comm. Math. Phys. '99

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Not complete yet. However  $u_k = \bar{u}_k + O(1)$ ,  $\bar{u}_k(r) = \int_{\partial B_r(x_k)} u_k$ 

In collaboration with M. Lucia: optimal sup + inf inequality (i.e. with  $C_1 = \frac{1}{n-1}$ ) when  $\sup_{k \in \mathbb{N}} \operatorname{osc}_{\partial B_{\delta}(x_k)} u_k < +\infty$ Not complete yet. However  $u_k = \bar{u}_k + O(1)$ ,  $\bar{u}_k(r) = \int_{\partial B_r(x_k)} u_k$ 

When  $n = 2 \ \bar{u}_k$  satisfies  $-\Delta \bar{u}_k = \int_{\partial B_r(x_k)} V_k e^{u_k}$  in  $B_{\delta}(x_k)$ .

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 $-\Delta_n v_k = \oint_{\partial B_r(x_k)} V_k e^{u_k} \text{ in } B_\delta(x_k), \quad v_k = \bar{u}_k(\delta) \text{ on } \partial B_\delta(x_k)$ 

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<u>Fundamental fact</u>:  $u_k - v_k = O(1)$ 

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<u>Fundamental fact</u>:  $u_k - v_k = O(1)$ <u>Crucial property</u>:  $-\Delta_n u_k + \Delta_n v_k = V_k e^{u_k} - \int_{\partial B_r(x_k)} V_k e^{u_k}$ 

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- optimal sup + inf inequality without boundary control
- more precise asymptotic expansion of  $u_k$  towards the case  $\lambda \in c_n \omega_n \mathbb{N}$
- blow-up solutions as  $\lambda \to c_n \omega_n \mathbb{N}$ . Via Lyapunov-Schmidt reduction: Pistoia, Premoselli, Vétois, ...
- multiplicity results via Morse theory

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# Thanks for your attention

P. Esposito May 16, 2024

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